

(8) The Addition of Angular Momenta

52

① Tensor products and direct sums

$$[\otimes]$$

$$[\oplus]$$

- They are to combine two Hilbert spaces into one.
(vector)

a. tensor products.

- For a vector $\vec{v} \in V$ and the other $\vec{w} \in W$

$$\rightarrow \vec{v} \otimes \vec{w} \in V \otimes W.$$

ex. a system of two spin $-\frac{1}{2}$ particles.

the state ket of the system:

$$\begin{aligned} |\Psi\rangle &= C_{\uparrow\uparrow} |\uparrow\rangle_1 \otimes |\uparrow\rangle_2 + C_{\uparrow\downarrow} |\uparrow\rangle_1 \otimes |\downarrow\rangle_2 \\ &\quad + C_{\downarrow\uparrow} |\downarrow\rangle_1 \otimes |\uparrow\rangle_2 + C_{\downarrow\downarrow} |\downarrow\rangle_1 \otimes |\downarrow\rangle_2 \end{aligned}$$

- Addition of two operators that are in different H-spaces.

$$O_v \in V, \quad O_w \in W$$

$$\Rightarrow O_v + O_w \equiv O_v \otimes I_w + I_v \otimes O_w$$

- Operator - ket multiplication
(matrix) (vector)

$$\rightarrow (S_1 + S_2) |\Psi\rangle = C_{\uparrow\uparrow} \left[(S_1 |\uparrow\rangle_1) \otimes |\uparrow\rangle_2 + |\uparrow\rangle_1 \otimes (S_2 |\uparrow\rangle_2) \right] + \dots$$

Rule $\uparrow\uparrow\uparrow$

$$\begin{aligned} &(O_v \otimes I_w + I_v \otimes O_w) \cdot (|v\rangle \otimes |w\rangle) \\ &= (O_v |v\rangle) \otimes (I_w |w\rangle) + (I_v |v\rangle) \otimes (O_w |w\rangle) \end{aligned}$$

- a. base kets and dimension of the combined space.

$$\begin{array}{ccc}
 \begin{array}{c} V \\ \{ |v_i\rangle \} \\ \dim[V] \end{array} & \otimes & \begin{array}{c} W \\ \{ |w_j\rangle \} \\ \dim[W] \end{array} \rightarrow \begin{array}{c} V \otimes W \\ \{ |v_i\rangle \otimes |w_j\rangle \} \\ \dim[V \otimes W] \\ = \dim[V] \cdot \dim[W] \end{array}
 \end{array}$$

- b. Direct sums.

$$\begin{array}{ccc}
 \begin{array}{c} V \\ \{ |v_i\rangle \} \\ \dim[V] \end{array} & \oplus & \begin{array}{c} W \\ \{ |w_j\rangle \} \\ \dim[W] \end{array} \rightarrow \begin{array}{c} V \oplus W \\ \{ |v_i\rangle \dots |w_j\rangle \} \\ \dim[V \oplus W] \\ = \dim[V] + \dim[W] \end{array}
 \end{array}$$

ex. matrix representation of $\mathcal{U}(R)$

$$\Rightarrow \mathcal{U}_{\text{min}}^{(j)}(R)$$

$$\left(\begin{array}{c} \boxed{j=1} \\ \boxed{j=2} \\ \vdots \end{array} \right) \Leftarrow V_{j=1} \oplus V_{j=2} \oplus \dots$$

"Block - diagonal".

We have already seen the tensor product many times. 74

$$\rightarrow |x, y, z\rangle = |x\rangle \otimes |y\rangle \otimes |z\rangle$$

$$H = \frac{\vec{p}^2}{2m} = \frac{1}{2m} \left((\tilde{p}_x \otimes I_y \otimes I_z)^2 + (I_x \otimes \tilde{p}_y \otimes I_z)^2 + (I_x \otimes I_y \otimes \tilde{p}_z)^2 \right)$$

C. notations.

1. DO NOT mix the index ordering.

$$\text{ex. } \cancel{| \uparrow \rangle_1 \otimes | \downarrow \rangle_2 + | \uparrow \rangle_2 \otimes | \downarrow \rangle_1} \quad (\text{X})$$

$$| \uparrow \rangle_1 \otimes | \downarrow \rangle_2 + | \downarrow \rangle_1 \otimes | \uparrow \rangle_2 \quad (\text{O})$$

2. $(\otimes)_{\text{index}}$ are often omitted for brevity.

$$\text{ex. } | \uparrow \rangle_1 \otimes | \downarrow \rangle_2 + | \downarrow \rangle_1 \otimes | \uparrow \rangle_2 \equiv | \uparrow \downarrow \rangle + | \downarrow \uparrow \rangle$$

: assuming that index ordering (1, 2) is fixed.

② Simple examples of Angular-Momentum addition

• base ~~ket~~ of a spin- $\frac{1}{2}$ particle in space.

$$|x; \uparrow\rangle \equiv |x\rangle \otimes | \uparrow \rangle, \quad |x; \downarrow\rangle \equiv |x\rangle \otimes | \downarrow \rangle$$

• Wave function:

$$\langle x; \uparrow | \alpha \rangle \equiv \psi_{\uparrow}(x), \quad \langle x; \downarrow | \alpha \rangle \equiv \psi_{\downarrow}(x).$$

$$\text{or } \Psi(x) \equiv \begin{pmatrix} \psi_{\uparrow}(x) \\ \psi_{\downarrow}(x) \end{pmatrix}$$

• L-S coupling. (spin-orbit interaction)

$$\vec{J} = \vec{L} + \vec{S} \quad (\vec{L} \cdot \vec{S} = \vec{J}^2 - \vec{L}^2 - \vec{S}^2)$$

$$\equiv \vec{L} \otimes I + I \otimes \vec{S} \quad \text{also. } \mathcal{H}(R) = \mathcal{H}_{(\vec{L})}^{(\text{orb})}(R) \otimes \mathcal{H}_{(\vec{S})}^{(\text{spin})}(R)$$

• Addition of $l=1$, $s=\frac{1}{2}$: $\vec{J} = \vec{L} + \vec{S}$

55

- base vectors (kets)

$$|s, m_s\rangle : |\frac{1}{2}, \frac{1}{2}\rangle \doteq \begin{pmatrix} 1 \\ 0 \end{pmatrix}, |\frac{1}{2}, -\frac{1}{2}\rangle \doteq \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

$$|l, m_l\rangle : |1, 1\rangle \doteq \begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix}, |1, 0\rangle \doteq \begin{pmatrix} 0 \\ 1 \\ 0 \end{pmatrix}, |1, -1\rangle \doteq \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix}$$

- Angular momentum operators

$$s=\frac{1}{2} : S_+ \doteq \hbar \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix}, S_- \doteq \hbar \begin{pmatrix} 0 & 0 \\ 1 & 0 \end{pmatrix}, S_z \doteq \frac{\hbar}{2} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

$$l=1 : L_+ \doteq \hbar \begin{pmatrix} 0 & \sqrt{2} & 0 \\ 0 & 0 & \sqrt{2} \\ 0 & 0 & 0 \end{pmatrix}, L_- \doteq \hbar \begin{pmatrix} 0 & 0 & 0 \\ \sqrt{2} & 0 & 0 \\ 0 & \sqrt{2} & 0 \end{pmatrix}, L_z \doteq \hbar \begin{pmatrix} 1 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & -1 \end{pmatrix}$$

↓ in the form of the tensor product $(l=1) \otimes (s=\frac{1}{2})$

$$L_+ \otimes I = \hbar \begin{pmatrix} \overset{m_l=1}{0} & \overset{m_l=0}{\sqrt{2}} & \overset{m_l=-1}{0} & 0 \\ 0 & 0 & \sqrt{2} & 0 \\ 0 & 0 & 0 & \sqrt{2} \\ 0 & 0 & 0 & 0 \end{pmatrix}$$

$$L_- \otimes I = \hbar \begin{pmatrix} 0 & 0 & 0 \\ \sqrt{2} & 0 & 0 \\ 0 & \sqrt{2} & 0 \\ 0 & 0 & 0 \end{pmatrix}$$

$$I \otimes S_+ = \hbar \begin{pmatrix} 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix}$$

$$I \otimes S_- = \hbar \begin{pmatrix} 0 & 0 & 0 & 0 \\ 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}$$

$$L_z \otimes I = \hbar \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix}$$

$$I \otimes S_z = \frac{\hbar}{2} \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix}$$

$$J_+ = L_+ \otimes I + I \otimes S_+ , \quad J_- = L_- \otimes I + I \otimes S_-$$

$$= \frac{1}{\hbar} \begin{pmatrix} 0 & 1 & \sqrt{2} & 0 \\ 0 & 0 & 0 & \sqrt{2} \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix}$$

$$= \frac{1}{\hbar} \begin{pmatrix} 0 & 0 & 0 & 0 \\ 1 & 0 & 0 & 0 \\ \sqrt{2} & 0 & 0 & 0 \\ 0 & \sqrt{2} & 1 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix}$$

$$J_z = L_z \otimes I + I \otimes S_z$$

$$= \frac{1}{\hbar} \begin{pmatrix} \frac{3}{2} & & & \\ & \frac{1}{2} & & \\ & & -\frac{1}{2} & \\ & & & -\frac{3}{2} \end{pmatrix}$$

see if

→

$$[J_i, J_j] = i\hbar \epsilon_{ijk} J_k$$

holds.

Yes!

eigenkets.

$|j, m\rangle$

→

It's obvious that

$$|\frac{3}{2}, \frac{3}{2}\rangle = \begin{pmatrix} 1 \\ 0 \\ 0 \\ 0 \end{pmatrix} \quad \begin{matrix} m_z = \frac{3}{2} \\ m_z = 1 \\ m_z = 0 \\ m_z = -1 \end{matrix} \quad , \quad |\frac{3}{2}, -\frac{3}{2}\rangle = \begin{pmatrix} 0 \\ 0 \\ 0 \\ 1 \end{pmatrix}$$

$$* \text{ check if } \begin{cases} J_+ |\frac{3}{2}, \frac{3}{2}\rangle = 0 \\ J_- |\frac{3}{2}, -\frac{3}{2}\rangle = 0 \end{cases}$$

Lowering with J_- :

$$J_- |\frac{3}{2}, \frac{3}{2}\rangle = \hbar \begin{pmatrix} 0 \\ 1 \\ 0 \\ 0 \end{pmatrix} = \hbar \sqrt{3} |\frac{3}{2}, \frac{1}{2}\rangle$$

--- (*)

$$\uparrow \text{ From } \sqrt{(j+m)(j-m+1)} \hbar |j, m-1\rangle = J_- |j, m\rangle$$

Raising with J_+ :

$$J_+ |\frac{3}{2}, -\frac{3}{2}\rangle = \hbar \begin{pmatrix} 0 \\ 0 \\ \sqrt{2} \\ 1 \end{pmatrix} = \hbar \sqrt{3} |\frac{3}{2}, -\frac{1}{2}\rangle$$

--- (*)

$$\uparrow \text{ From } \sqrt{(j-m)(j+m+1)} \hbar |j, m+1\rangle = J_+ |j, m\rangle$$

also, by just using operators and kets.

$$J_- \left| \frac{3}{2}, \frac{3}{2} \right\rangle = (L_- \otimes I + I \otimes S_-) \left(\overset{L \ m_L}{|1,1\rangle} \otimes \overset{S \ m_S}{\left| \frac{1}{2}, \frac{1}{2} \right\rangle} \right) \\ = \hbar\sqrt{2} \left| 1,0 \right\rangle \otimes \left| \frac{1}{2}, \frac{1}{2} \right\rangle + \hbar \left| 1,1 \right\rangle \otimes \left| \frac{1}{2}, -\frac{1}{2} \right\rangle = \hbar\sqrt{3} \left| \frac{3}{2}, \frac{1}{2} \right\rangle$$

$$\Rightarrow \left| \frac{3}{2}, \frac{1}{2} \right\rangle = \sqrt{\frac{1}{3}} \left| 1,0 \right\rangle \otimes \left| \frac{1}{2}, \frac{1}{2} \right\rangle + \sqrt{\frac{2}{3}} \left| 1,1 \right\rangle \otimes \left| \frac{1}{2}, -\frac{1}{2} \right\rangle \\ = \begin{pmatrix} 0 \\ 0 \\ \frac{1}{\sqrt{3}} \\ \frac{\sqrt{2}}{\sqrt{3}} \\ 0 \\ 0 \end{pmatrix}$$

$$J_+ \left| \frac{3}{2}, -\frac{3}{2} \right\rangle = (L_+ \otimes I + I \otimes S_+) \left(\left| 1,-1 \right\rangle \otimes \left| \frac{1}{2}, -\frac{1}{2} \right\rangle \right) \\ = \hbar \left| 1,-1 \right\rangle \otimes \left| \frac{1}{2}, \frac{1}{2} \right\rangle + \hbar\sqrt{2} \left| 1,0 \right\rangle \otimes \left| \frac{1}{2}, -\frac{1}{2} \right\rangle = \hbar\sqrt{3} \left| \frac{3}{2}, -\frac{1}{2} \right\rangle$$

$$\Rightarrow \left| \frac{3}{2}, -\frac{1}{2} \right\rangle = \frac{1}{\sqrt{3}} \left| 1,-1 \right\rangle \otimes \left| \frac{1}{2}, \frac{1}{2} \right\rangle + \sqrt{\frac{2}{3}} \left| 1,0 \right\rangle \otimes \left| \frac{1}{2}, -\frac{1}{2} \right\rangle \\ = \begin{pmatrix} 0 \\ 0 \\ \frac{1}{\sqrt{3}} \\ \frac{\sqrt{2}}{\sqrt{3}} \\ 0 \\ 0 \end{pmatrix}$$

by using the orthogonality & "convention",

we may write

$$\left| \frac{1}{2}, \frac{1}{2} \right\rangle = -\sqrt{\frac{1}{3}} \left| 1,0 \right\rangle \otimes \left| \frac{1}{2}, \frac{1}{2} \right\rangle + \sqrt{\frac{2}{3}} \left| 1,1 \right\rangle \otimes \left| \frac{1}{2}, -\frac{1}{2} \right\rangle$$

$$\left| \frac{1}{2}, -\frac{1}{2} \right\rangle = -\sqrt{\frac{2}{3}} \left| 1,-1 \right\rangle \otimes \left| \frac{1}{2}, \frac{1}{2} \right\rangle + \sqrt{\frac{1}{3}} \left| 1,0 \right\rangle \otimes \left| \frac{1}{2}, -\frac{1}{2} \right\rangle$$

↓

✓

⇒ "Clebsch-Gordan Coefficients"
☆☆☆

$$\begin{pmatrix} \left| \frac{3}{2}, \frac{1}{2} \right\rangle \\ \left| \frac{1}{2}, +\frac{1}{2} \right\rangle \end{pmatrix} = \begin{pmatrix} \sqrt{\frac{1}{3}} & \sqrt{\frac{1}{3}} \\ -\sqrt{\frac{1}{3}} & \sqrt{\frac{2}{3}} \end{pmatrix} \begin{pmatrix} \left| 1,0 \right\rangle \otimes \left| \frac{1}{2}, \frac{1}{2} \right\rangle \\ \left| 1,1 \right\rangle \otimes \left| \frac{1}{2}, -\frac{1}{2} \right\rangle \end{pmatrix}$$

orthogonal matrix.

So, we may construct a unitary transformation.

58

from $\{|l, m_l\rangle \otimes |s, m_s\rangle\}$ to $\{|j, m\rangle\}$ as

$$1 \otimes \frac{1}{2} = \frac{3}{2} \oplus \frac{1}{2}$$

$$: 3 \times 2 = 4 + 2$$

$$U(|l, m_l\rangle \otimes |s, m_s\rangle) = |j, m\rangle$$

$$\Rightarrow U = \begin{pmatrix} 1 & 0 & 0 & 0 & 0 & 0 \\ 0 & \sqrt{1/3} & \sqrt{2/3} & 0 & 0 & 0 \\ 0 & 0 & 0 & \sqrt{1/3} & 1/\sqrt{3} & 0 \\ 0 & 0 & 0 & 0 & 0 & 1 \\ 0 & \sqrt{2/3} & -1/\sqrt{3} & 0 & 0 & 0 \\ 0 & 0 & 0 & 1/\sqrt{3} & -\sqrt{2/3} & 0 \end{pmatrix}$$

"multiplet"

(cf.) triplet
singlet

in $\frac{1}{2} \otimes \frac{1}{2}$

Verify:

$$U J_z U^\dagger = \hbar \begin{pmatrix} \frac{3}{2} & & & & & \\ & \frac{1}{2} & & & & \\ & & -\frac{1}{2} & & & \\ & & & -\frac{3}{2} & & \\ & & & & \frac{1}{2} & \\ & & & & & -\frac{1}{2} \end{pmatrix}$$

* How to read out the table of CG coefficients

☆☆☆☆

Just Google it!
you'll find the table.

$$1 \times \frac{1}{2}$$

	$\frac{3}{2}$	$\frac{1}{2}$
$+\frac{1}{2}$	1	$+\frac{1}{2}$
0	$+\frac{1}{2}$	$-\frac{1}{2}$

$\frac{3}{2}, +\frac{1}{2}$

$1, 1 \otimes \frac{1}{2}, -\frac{1}{2}$

square-root is omitted

$$\sqrt{2/3}$$

J	M	...
m_1	m_2	
m_1	m_2	CG coeff.
\vdots	\vdots	

$$| \frac{3}{2}, \frac{1}{2} \rangle = \frac{1}{\sqrt{3}} | 1, 1 \rangle \otimes | \frac{1}{2}, -\frac{1}{2} \rangle + \sqrt{\frac{2}{3}} | 1, 0 \rangle \otimes | \frac{1}{2}, \frac{1}{2} \rangle$$

$$\text{also, } | 1, 1 \rangle \otimes | \frac{1}{2}, -\frac{1}{2} \rangle = \frac{1}{\sqrt{3}} | \frac{3}{2}, \frac{1}{2} \rangle + \sqrt{\frac{2}{3}} | \frac{1}{2}, \frac{1}{2} \rangle$$

There're also computation codes to calculate CG coeffs.

In python, c/c++, ...